

From multipartite entanglement to TQFTs

Joint w/ Pavel Putrov & Abhijit Gadde
entanglement measures

$$\mathcal{H} = \mathcal{H}_1 \otimes \mathcal{H}_2 \otimes \dots \otimes \mathcal{H}_p$$

$p = \#$ of parties

$$f_{\alpha_i} \quad i = 1, \dots, p$$
$$\alpha_i = 1, \dots, \dim \mathcal{H}_i$$

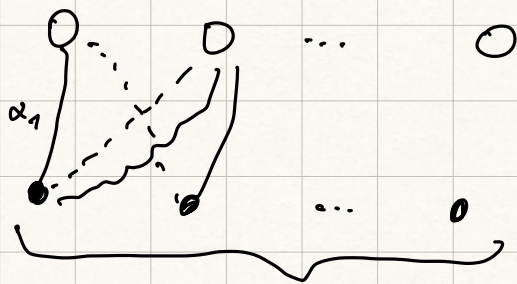
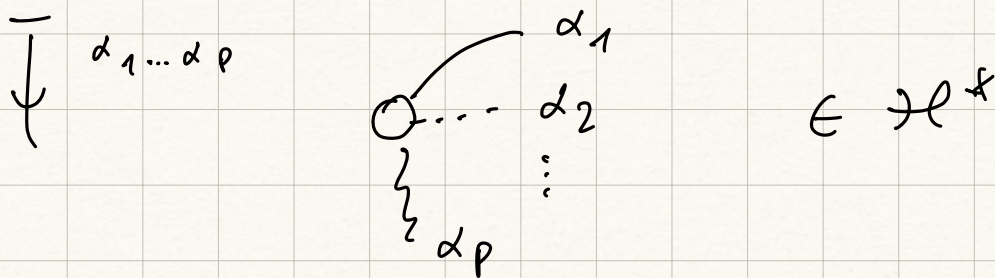
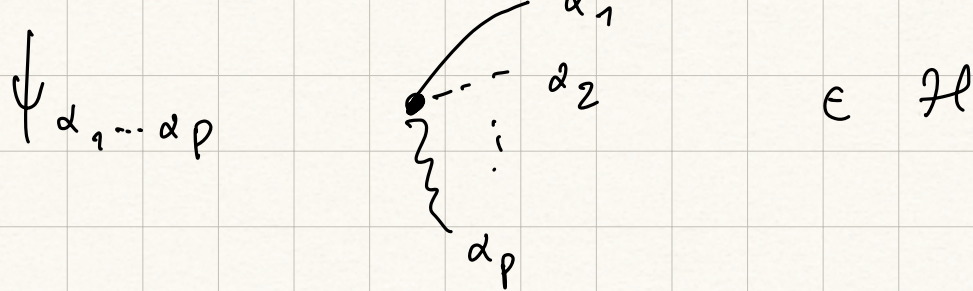
$$\Psi = \sum_{\alpha_i} \left(\psi_{\alpha_1 \dots \alpha_p} f_{\alpha_1} \otimes \dots \otimes f_{\alpha_p} \right)$$

to measure entanglement

into an \mathbb{Z} that is invariant under
local unitary transformations

$$U = U_1 \otimes U_2 \otimes \dots \otimes U_p$$

$$U_i \in \mathcal{U}(\mathcal{H}_i)$$



$R = \#$ OF REPLICAS

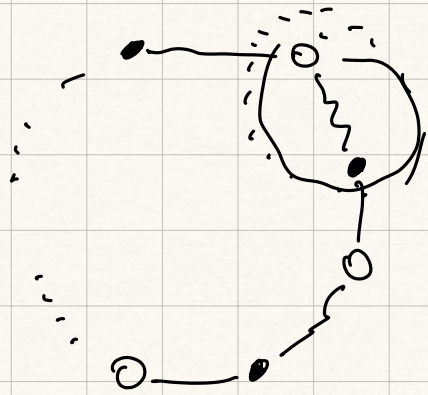
$\leadsto \sum_P (\Psi)$ \prod_P bipartite graph
 "entanglement measure"

example

$p = 2$

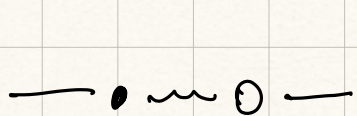
$\mathcal{H} = \mathcal{H}_1 \otimes \mathcal{H}_2$



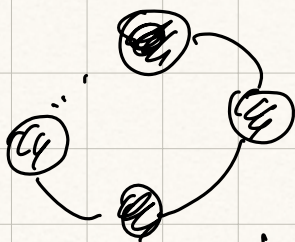


$$R = \# \bullet$$

$\left\{ \begin{array}{l} R\text{-th Rényi} \\ \text{entropy} \end{array} \right.$



$$\rho_{\alpha_1}^{\alpha_1'} = \sum_{\alpha_2} \psi_{\alpha_1 \alpha_2} \overline{\psi}^{\alpha_1' \alpha_2}$$



$$\text{Tr } \rho^R$$

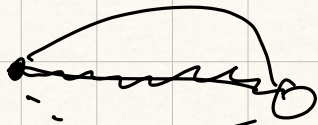
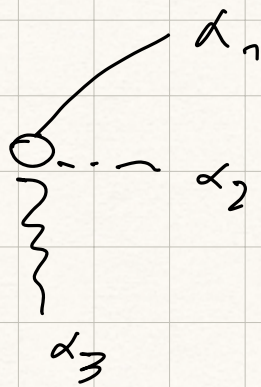
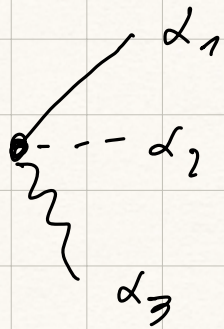
$$\lim_{R \rightarrow 1} \frac{1}{1-R} \log \text{tr } \rho^R$$

\leadsto von Neumann entropy \mathcal{J}

Multipartite entanglement

one could impose / ask further
 properties (vanish on prod. state
 monotonicity ...)

$$p = 3$$



RMC

Here

we
have
generalized

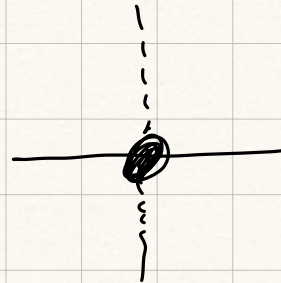
$\cup N$

type

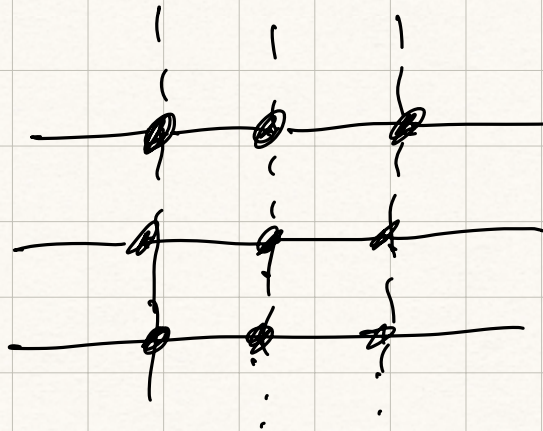
p -point
entropies



\rightsquigarrow



N



reduced
density

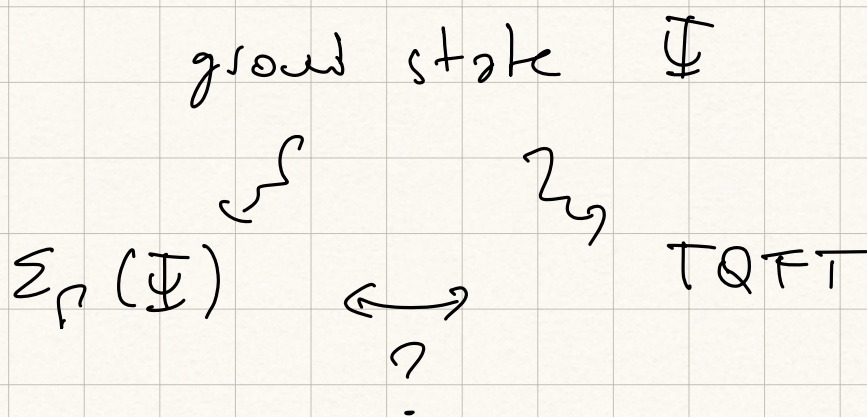
"tensor"

$$\rightsquigarrow Z_N^{(3)}$$

$$S^{(3)} = \frac{1}{1-N} \lim_{N \rightarrow 1} \ln \ln Z_N^{(3)}$$

good characterization of genuine

N -point entanglement \rightarrow OPEN?



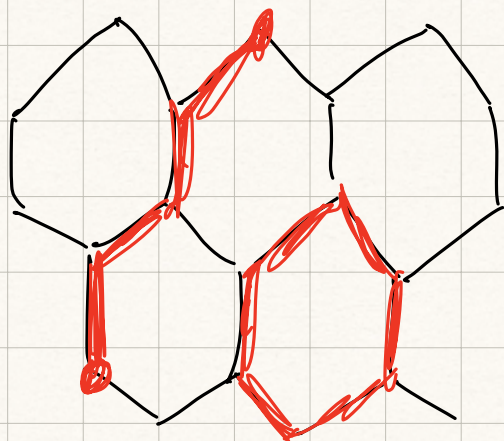
gapped lattice models

\rightarrow g.s. are described by TQFT

string-net model Levin-Wen

\leadsto (and of Turaev-Viro

ground states are known



string net
configurations
|||

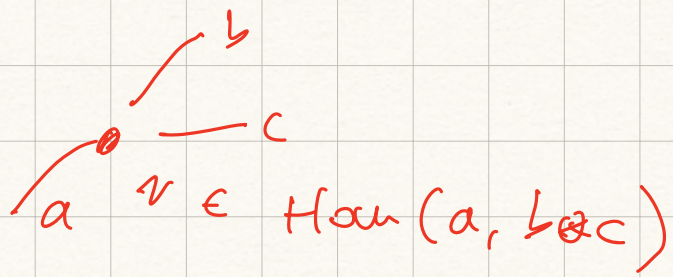
graphical
calculus
of spherical
fusion CAT \mathcal{C}

in the IR

→ Turaev
Viro

TQFT

for \mathcal{E}



* Kirillov Jr

& Bolson

~ 2010

ground state Ψ_{LW}

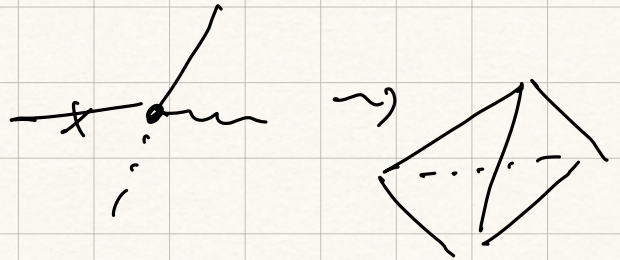
$$\textcircled{\star} \left| \log \sum_{\mathcal{F}_4} (\Psi_{LW}) = \log Z_{TV} (M_{\mathcal{F}_4}^{(\mathcal{E})}) + O(e^{-cL}) \right.$$

Graphically
encoded 3-manifolds

generalizes

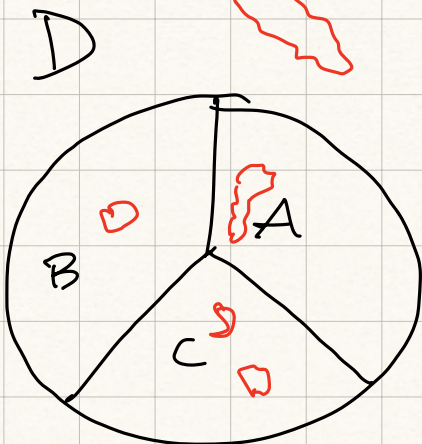
Kitaeu-Preisk-U

Levin-Wen / 2006



planar
hexagonal
lattice

divided into
4 (reg) regions

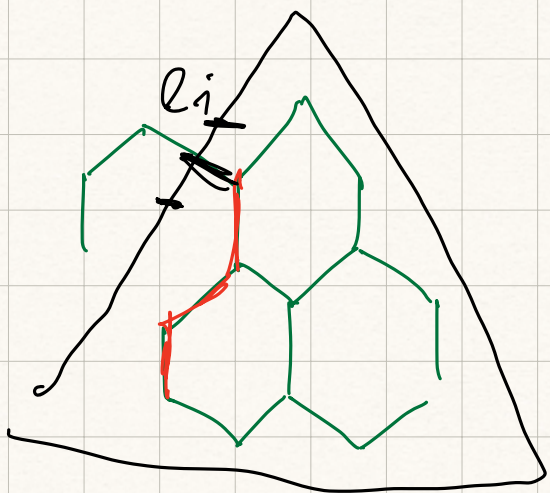


~) 4-partite entanglement
regions

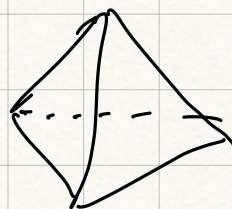
$$\mathcal{H} = \mathcal{H}_A \otimes \mathcal{H}_B \otimes \mathcal{H}_C \otimes \mathcal{H}_D$$

strings don't go to ∞

\hookrightarrow topologically trivial on S^2



S
yours



$$l_i \in \text{tr}(\mathcal{E})$$

$$\mathcal{H} = \text{Hom}_{\mathcal{E}}(1, l_1 \otimes \dots \otimes l_n)$$

Ψ
CW

can construct it
explicitly

using Kac-Moody
work

we can show (\star)
provided

$$\sum_{l \in \text{tr}(\mathcal{E})^L} N_{l_1 \dots l_L} \prod_{i=1}^L d_{e_i}^u$$

$$= \frac{\left(\sum_{e \in \text{irr}(e)} d_e^{u+1} \right)^L d_{\vec{e}}}{(1 + O(k^{-c}))}$$

$$\mathbb{D}^2 = \left(\sum_{e \in \text{irr}(e)} d_e^2 \right)$$

$L \gg 1$

$c > 0$

$$N_{e_1 \dots e_L}^{\vec{e}} = \dim \text{Hom}_{\mathbb{C}}(e_1 \otimes \dots \otimes e_L, \vec{e})$$

$n=1$ is well-known

$$\sum_{e \in \text{irr}(e)} N_{e_1 \dots e_L}^{\vec{e}} \prod_{i=1}^L d_{e_i} = \mathbb{D}^{2L-2} d_{\vec{e}}$$

$n > 1$ Povel can prove it when e is an MTC

In general we expect

d -dim \mathbb{C}

TQFT

on M_d^d
 Γ_{d+1}

\rightarrow

$(d+1)$ -particle

entangled

states

\sum_{Ψ_0}

glueing Δ^{d+1}
simplexes

$d+1$

Questions

1.) what do we need to
completely classify
a TQFT?

having all partition
functions is it enough?

2.) what is special of
the $\mathbb{R} \rightarrow$ manifolds

(genuine multiple
entanglement?)

3.) when we classify phases
out orders

→ should we ask we
retain info about
entanglement?

Graph-encoded manifolds

